

Non-Linear model to simulate large amplitude motions: application to wave energy conversion.

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INTRODUCTION

In this paper we present a time-domain non-linear potential-flow model to predict large-amplitude motions of the SEAREV device [1] in a prescribed or arbitrary seaway. The body is a floating device, completely closed, with an internal moving mass. Under the action of the waves, the floating hull and the internal moving mass start to move, each one with its own motion. The relative motion between the floating body and the moving mass is used to drive a generator via an hydraulic device, the Power Take Off (PTO). In the model, the floating body is three-dimensional, undergoing arbitrary six degree-of-freedom motions under the action of incoming waves. The dynamic and static components of the non-linear Froude-Krylov forces are computed on the instantaneous wetted surface under the incident wave profile. A linear time-domain model is applied to compute the first-order of the radiation and diffraction forces. These forces are developed to the second-order by using Taylor expansion. The quadratic component of the Bernoulli's equation is added by convolution. The test case of the vertical cylinder is studied to compare the mean surge force with analytic solution. Comparisons with experimental results are also made for the SEAREV device in regular seas.

DESCRIPTION OF THE MODEL

Coordinate system convention.

The floating body, is assumed to have two vertical planes of symmetry. We assume that the axis of the internal cylinder is perpendicular to the main symmetry plane of the floating body. We also assume that the direction of propagation of the waves can be out of this plane.

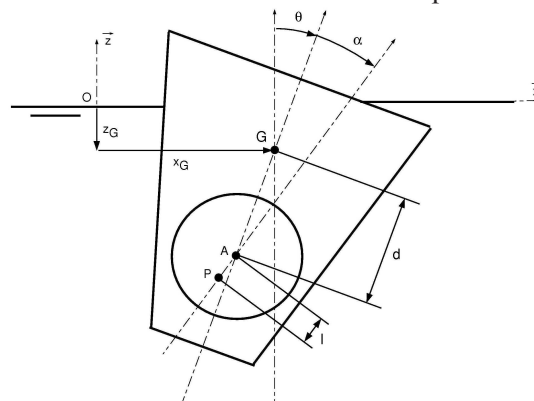


Figure 1 : Notations.

Let x_G be the surge motion, y_G be the sway motion and z_G the heave motion of the center of gravity G of the floating body. Let φ the roll motion along the x -axis, θ the pitch motion along the y -axis, ψ the yaw motion along the z -axis and α the relative motion between the floating body and the inner pendulum. Note that the PTO system functions due to α . Let m_b be the floating body mass, I_b its inertia, m_p the inner pendulum mass and I_{pb} its inertia. Let $l=AP$ be the distance between the center of rotation A of the inner pendulum and its own center of gravity P and $d=GA$ the distance between the center of gravity of the floating body and the center of rotation of the pendulum.

Equations of motion

The SEAREV wave energy converter is a seven degrees-of-freedom device. Let $\mathbf{X} = (x_G \ y_G \ z_G \ \varphi \ \theta \ \psi \ \alpha)^t$ be the position vector. In the time domain, the equations of motion are:

$$\mathbf{M}(\mathbf{X}, \dot{\mathbf{X}})\ddot{\mathbf{X}} = \mathbf{F}_c(\mathbf{X}, \dot{\mathbf{X}}) + \mathbf{F}_{\text{PTO}} + \mathbf{F}_H$$

where

- \mathbf{M} is the 7x7 proper mass matrix of the device. This matrix contains coupling terms between the pendulum and the floating body and depends on the position vector and the velocity vector, as we did not assumed linearity. The calculation of \mathbf{M} is complex and is not the main point of this paper. Details and expression of \mathbf{M} as a function of the mechanical parameters of the device can be found in [2].
- \mathbf{F}_C is the Coriolis force. The expression of this force is given in [2].
- \mathbf{F}_{PTO} represents the action of the Power-Take-Off. It is modelled here as a linear damper, with B_{PTO} being the damping coefficient.
- \mathbf{F}_H represents the force due to the fluid-structure interactions.

Fluid-structure interactions

The following usual assumptions are made: the fluid is considered homogeneous, incompressible, inviscid and the flow irrotational. Surface tension is not taken into account, the depth is considered infinite and a linearized free surface and body boundary conditions are used. The fluid forces acting on the body can be non-linear with respect to certain motion variables, e.g the quadratic component of the Bernoulli’s equation, the nonlinear incident potential flow. They contain “geometric” non-linearities as in the case where the force are computed by integrating on the exact instantaneous position and shape on the immersed part of the body. The first-order force is calculated by a linear potential flow formulation whereas the second-order force is calculated by adding the quadratic term of Bernoulli’s equation and by expanding the first-order force to the second-order with the Taylor expansion

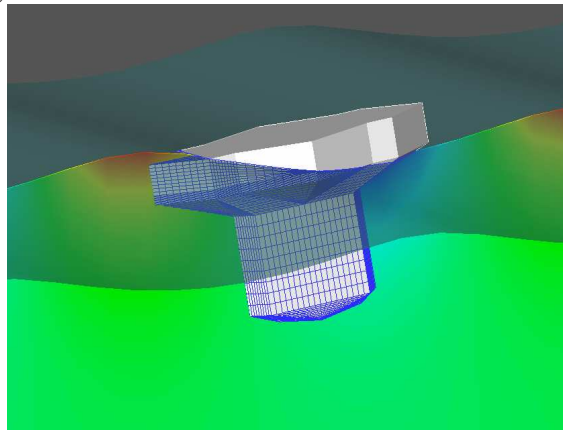


Figure 2 : 3D view of the SEAREV and its wetted surface.

Froude-Krylov forces

The integration of the incident wave pressure is performed over the instantaneous total wetted surface at each time step, taking into account the large motions of the ship and its exact intersection with the undisturbed wave surface. The Froude-Krylov force is expressed by:

$$f_{FK}(t) = \iint (p_S + p_D) n ds$$

Where the static and dynamic pressure are, respectively:

$$p_S = -\rho g z, \quad -\infty < z < \eta$$

$$p_D = -\rho \frac{\delta\phi_I}{\delta t}$$

The dynamic pressure is derived from the non-linear incident potential of a higher-order spectral method [8]. A robust geometry processing capability is essential for large amplitude motions. Consequently, an automatic remeshing routine was developed for this task. At each time step, the underwater geometry is represented by a number of panels. As the body moves, its new location and orientation is updated in the global coordinate system and the new waterline is found from the intersection with the instantaneous free surface. The underwater portion of each panel is then repanelized using the transfinite method [9] (Figure (2)).

Radiation forces

The linear radiation forces have been expressed as a convolution product according to the well-known Cummins’ decomposition:

$$F_{\text{rad}}(t) = -\mu_{\infty} \dot{Y} - \int_0^t K(t-\tau) \dot{X}(\tau) d\tau$$

It can first be transformed in order to remove the convolution product by using Prony's method. This method has been implemented by Daubisse [5] and further developed by Clément [4] for the computation of impulse response of radiation forces

Diffraction forces

Like the radiation forces, the diffraction forces are based here on linear theory in the time domain. The diffracted wave forces are computed as:

$$F_{\text{diff}}(t) = \int_{-\infty}^{\infty} K_7(t-\tau) \eta_0(\tau) d\tau$$

where K_7 is the impulse response function for the diffraction forces and η_0 is the free surface elevation of the incident wave at a given reference point.

Expansion to the second-order.

The expansion to the second order is realized in two steps. In the first step, the linear hydrodynamic force is developed, using Taylor expansion in order to obtain:

$$\begin{aligned} F_{2nd}(t) = & \iint_{C_0} \phi_t(t, M'_0) n'_0 dM'_0 \\ & + \iint_{C_0} \left[\partial x' \cdot \frac{\partial \phi_t(t, M'_0)}{\partial x'} + \partial y' \cdot \frac{\partial \phi_t(t, M'_0)}{\partial y'} + \partial z' \cdot \frac{\partial \phi_t(t, M'_0)}{\partial z'} \right] n'_0 dM'_0 \\ & + \iint_{C_0} \phi_t(t, M'_0) \partial n'_0 dM'_0 \end{aligned}$$

In the second step, the quadratic term of the Bernoulli's equation is taken into account by a convolution as the following:

$$F_{V^2/2}(t) = -\frac{\rho}{2} \iint_{C_0} |\nabla \phi \cdot \dot{X}|^2 n_0 dS - \frac{\rho}{2} \iint_{C_0} |\nabla \phi \cdot V_I|^2 n_0 dS$$

where C_0 is the mean wetted surface, V_I is the incident wave. Terms from second-order radiation and diffraction potentials were ignored in this study.

VALIDATION

We considered a truncated circular cylinder of radius 1 meter in infinite water depth. The mean drift forces on the cylinder are showed on the figure (3). The forces are normalized by $\rho g R_0 A^2$ where ρ is the water density, g is the gravitational acceleration, R_0 is the radius and A is the wave amplitude. K is the infinite depth wave number. The plot on the figure (3) shows surge forces computed by two different approaches; the analytic solution (Molin [6]) and the model presented in this paper. The figure (3) shows that the numerical model underestimates the mean surge force in the low wave number range.

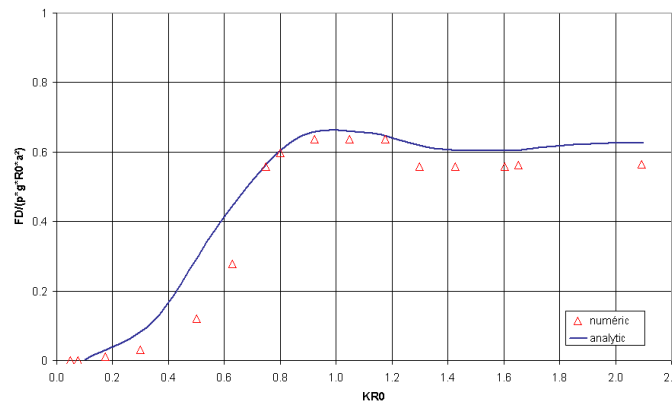


Figure 3 : Nonlinear mean surge forces on a cylinder.

Correlation between numerical and experimental results.

Experiments with 1/12-scale model of the SEAREV body were conducted in ECN's ocean engineering basin [7]. In order to compare the experimental results against the numerical model solutions, all the results were condensed in the form of transfer functions. Figure (5) shows a comparison between the non-linear formulation, a linear time domain formulation [2] and experimental results for a regular wave train with 0,083 meter of amplitude and with frequency is the range between 0,4 and 0,64 Hz. A fair agreement between theory and experiments is thus obtained. Figure (4) shows that a good prediction of the behaviour of the SEAREV device is obtained. The dynamic of the system is correctly recovered even if the motion amplitudes are slightly overestimated.

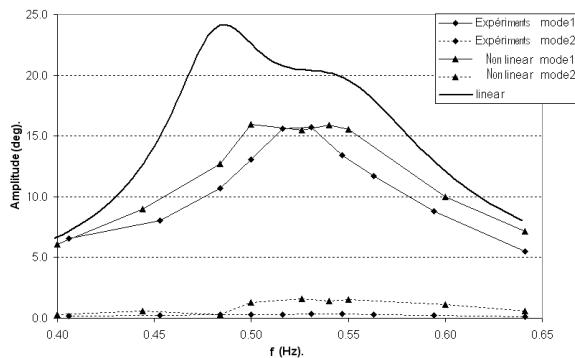


Figure 5: RAO for the pitch motion of the SEAREV device. Comparison between experiments and numerical simulation.

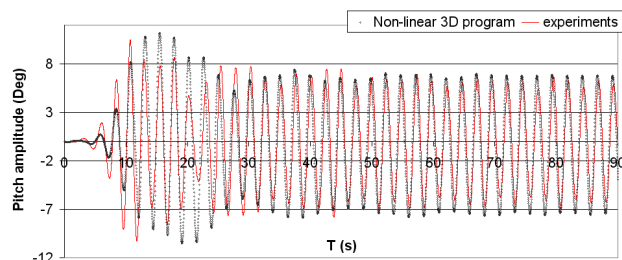


Figure 4: Pitch angle as a function of time. Comparison between experiments and numerical simulation.

CONCLUSION

We have presented in this paper a numerical model to determine the large amplitude motions of a floating wave energy converter subjected to incoming regular waves. The numerically motions of the body in regular waves, were compared with experimental results. The results predicted by the model agree reasonable well with experimental results.

Acknowledgments

This work is a result of research sponsored in part by ADEME, Agence de l'environnement et de la Maîtrise de l'Energie.

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